



RANDOM
GRAPHS
COMPLEX
NETWORKS

G. Panasenko

VJM - Univ de Lyon

ICJ - MODMAD

Equations on a graph
for the flows in thin
tube structures

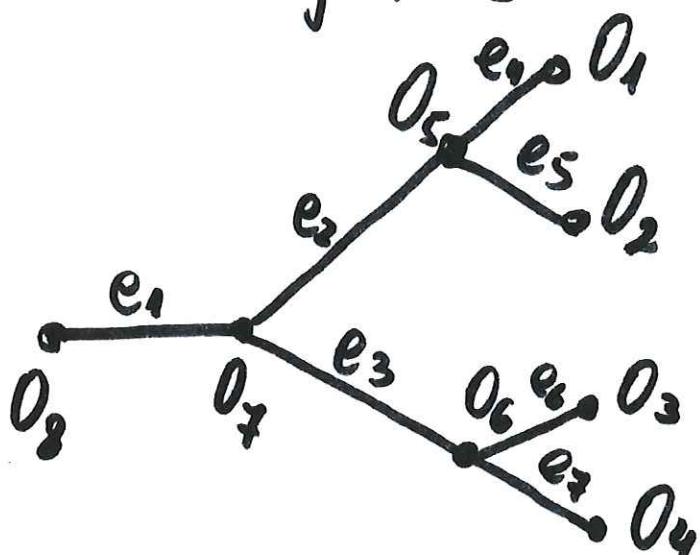
THIN STRUCTURE

1. Graph. $O_1, \dots, O_N \in \mathbb{R}^n$, $n \in \{2, 3\}$

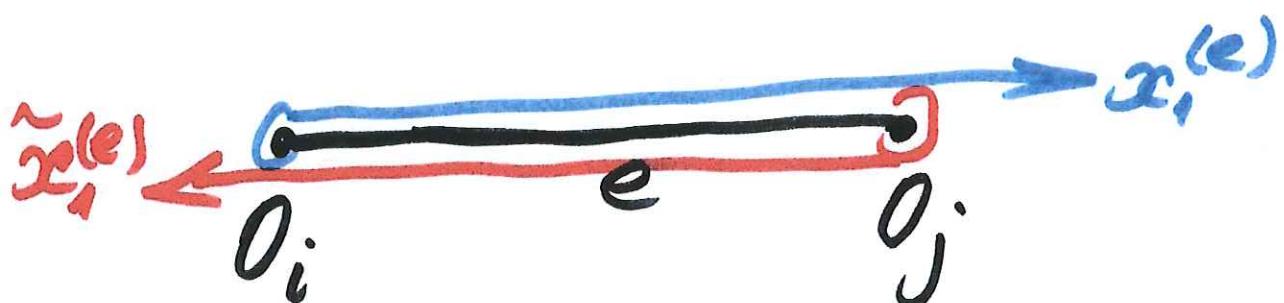
e_1, \dots, e_M closed segments

$$e_j = \overline{O_{i_j} O_{k_j}}, e_j \cap e_{j_2} \subseteq \{O_1, \dots, O_N\}$$

$$\mathcal{B} = \bigcup_{j=1}^M e_j$$

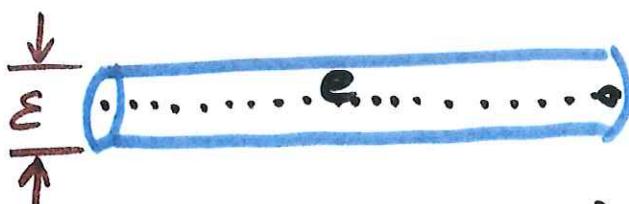


2. Local coordinate systems



3. Thin structure

$$e_j = e \quad \sigma^j = \sigma^{(e)}$$



$$B_\epsilon^{(e)} = \{x^{(e)} \mid x_i^{(e)} \in (0, |e|), \frac{x^{(e)}}{\epsilon} \in \sigma^{(e)}\}$$

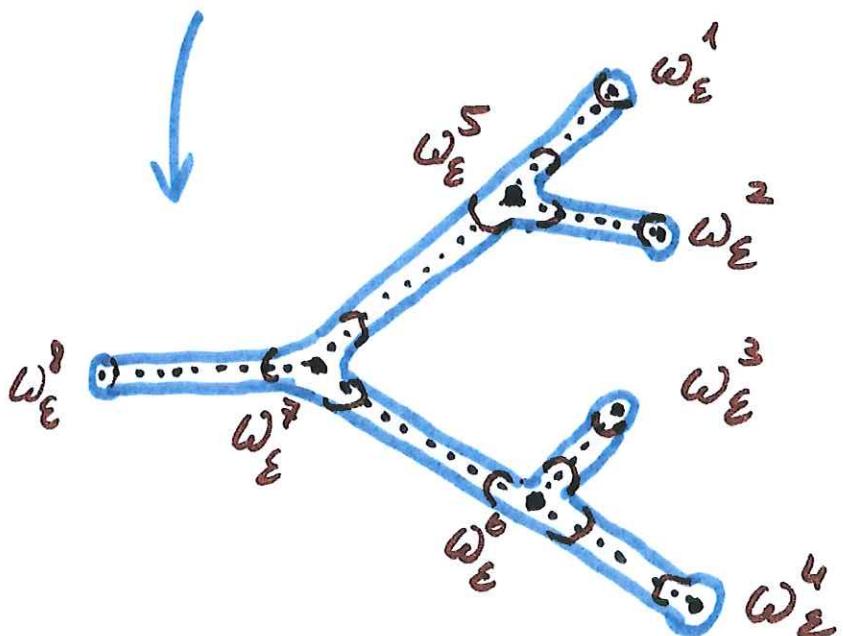
ω^i bounded domains
in \mathbb{R}^n , $i = 1, \dots, N$

$$\boxed{x^{(e)} \in \sigma_\epsilon^{(e)}}$$

$$\omega_\epsilon^i = \left\{ x \in \mathbb{R}^n \mid \frac{x - o_i}{\epsilon} \in \omega^i \right\}$$

$$B_\epsilon = \left(\bigcup_{j=1}^m B_\epsilon^{(e_j)} \right) \cup \left(\bigcup_{j=1}^N \omega_\epsilon^j \right)$$

connected
 $\partial B_\epsilon \in C^2$



$$\underline{\underline{\omega_i}}^E \quad \gamma_\varepsilon^i = \partial \omega_i^i \cap \partial B_\varepsilon \\ i = N_1 + 1, \dots, N : O_i \text{ vertices}$$

4. Navier - Stokes equations

$$\frac{1}{\varepsilon^2} \frac{\partial \vec{u}_\varepsilon}{\partial t} - \nu \Delta \vec{u}_\varepsilon + (\vec{u}_\varepsilon, \nabla) \vec{u}_\varepsilon + \nabla p = 0 \\ x \in B_\varepsilon, t \in \mathbb{R}$$

$$\operatorname{div} \vec{u}_\varepsilon = 0$$

$$\vec{u}_\varepsilon = \vec{g}_i \left(\frac{x - O_i}{\varepsilon}, t \right), x \in \gamma_\varepsilon^i \quad (1)$$

$$\vec{u}_\varepsilon = 0, x \in \partial B_\varepsilon \setminus Y_\varepsilon, Y_\varepsilon = \bigcup_{i=N_1+1}^N \gamma_\varepsilon^i$$

$$\vec{u}_\varepsilon(x, t+T) = u_\varepsilon(x, t)$$

\vec{g}_i : compact support functions on $\partial \omega_i$

T-periodic in time

$$\sum_i \int_{\gamma_\varepsilon^i} \vec{g}_i \cdot \vec{n} = 0$$

$$\varepsilon \ll 1$$

or

$$\frac{\partial \vec{u}_\varepsilon}{\partial t} - \nu \Delta \vec{u}_\varepsilon + (\vec{u}_\varepsilon, \nabla) \vec{u}_\varepsilon + \nabla p = 0$$

$x \in B_\varepsilon, t \in \mathbb{R},$

$$\operatorname{div} \vec{u}_\varepsilon = 0$$

$$\vec{u}_\varepsilon = \vec{g}_i \left(\frac{x - O_i}{\varepsilon}, \frac{t}{\varepsilon^2} \right), x \in \gamma_\varepsilon^i \quad (1_2)$$

$$\vec{u}_\varepsilon = 0, x \in \partial B_\varepsilon \setminus \gamma_\varepsilon$$

$$\vec{u}_\varepsilon(x, t + \varepsilon^2 T) = \vec{u}_\varepsilon(x, t)$$

or steady problem

$$-\nu \Delta \vec{u}_\varepsilon + (\vec{u}_\varepsilon, \nabla) \vec{u}_\varepsilon + \nabla p = 0$$

$x \in B_\varepsilon,$

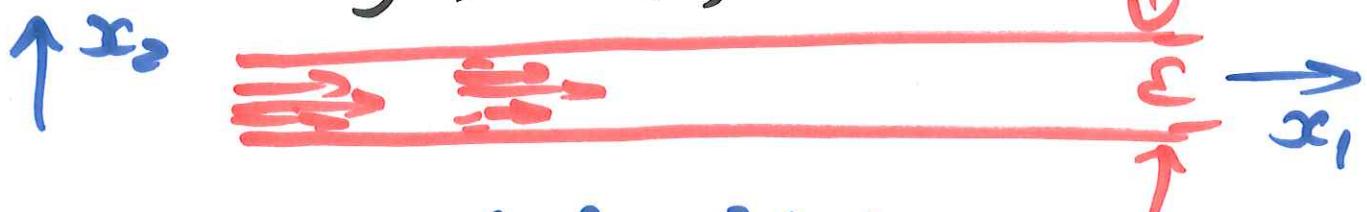
$$\operatorname{div} \vec{u}_\varepsilon = 0, \quad (1_3)$$

$$\vec{u}_\varepsilon = \vec{g}_i \left(\frac{x - O_i}{\varepsilon} \right), x \in \gamma_\varepsilon^i,$$

Complete as. expansion
is built.

1. For every edge the
Poiseuille flow is built

- steady, $n=2$,



$$u_1 = -a(x_2^2 - \frac{\epsilon^2}{2})/2y, u_2 = 0,$$

$$P = -a x_1 + b$$

- steady, $n=3$,

$$u_1(x_2, x_3) \Rightarrow a U_p, U_p : -\Delta_{\perp} U_p = 1$$

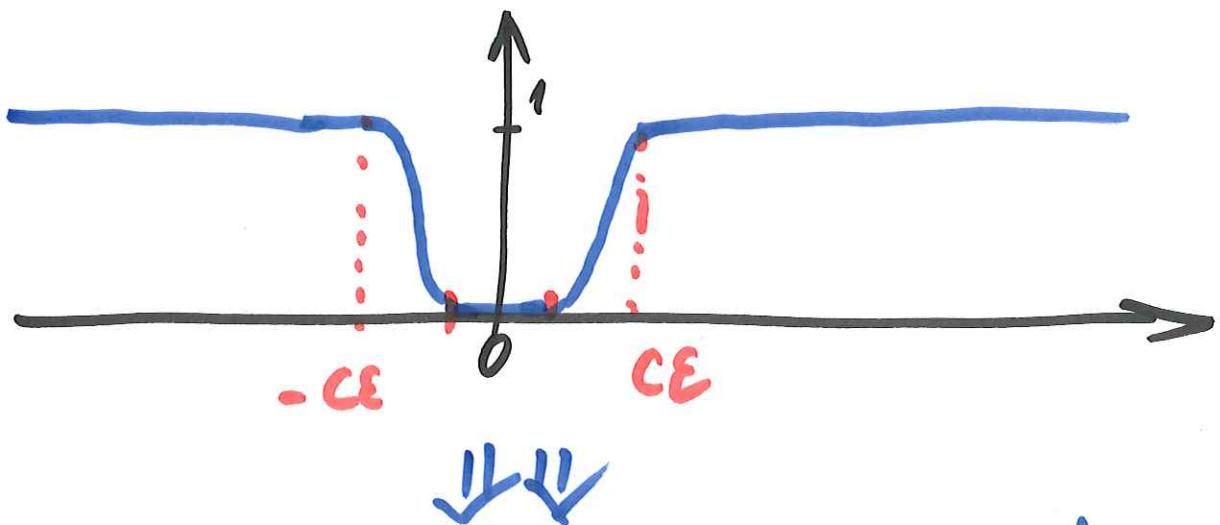
$$u_2 = u_3 = 0, \quad \boxed{on \overline{\Omega_\epsilon}, U_p|_{\partial\Omega_\epsilon} = 0}$$

$$P = -a x_1 + b$$

- non-steady : $u_1(x_2, x_3, t), u_2 = u_3 = 0,$

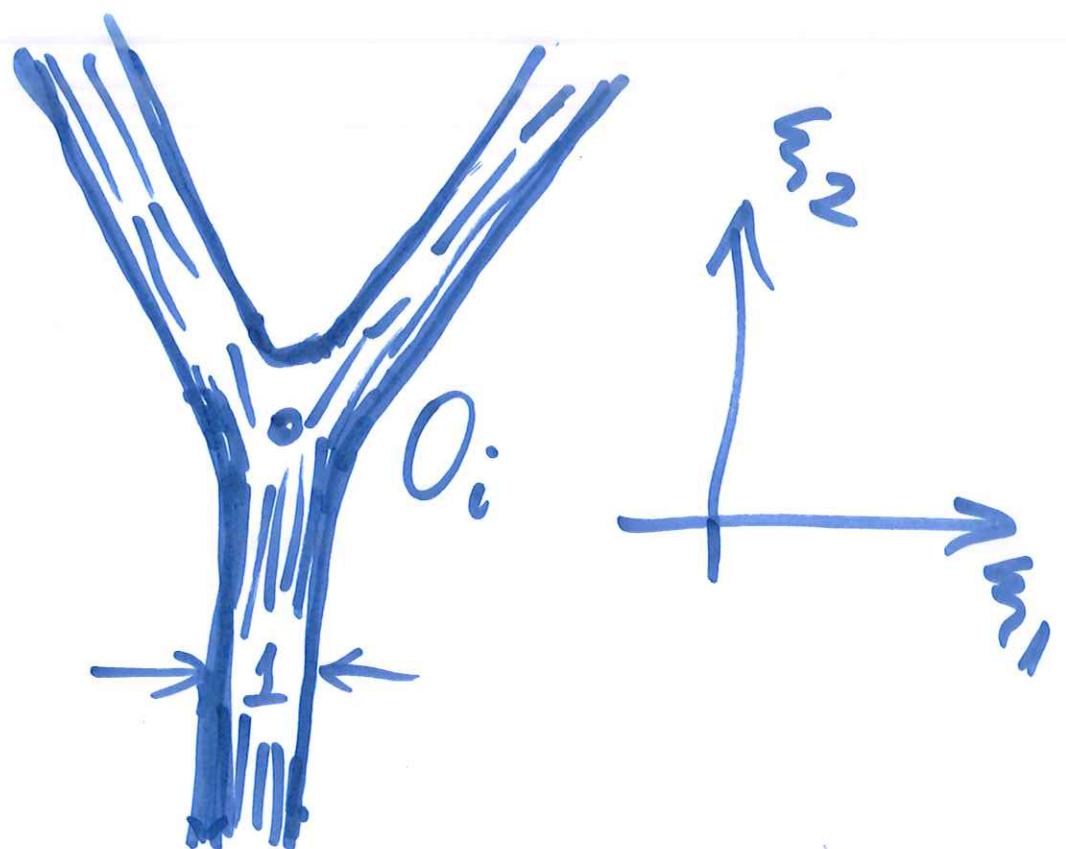
$$P = -a x_1 + b.$$

2. Each Poiseuille is multiplied by a smooth cut-off function vanishing at some neighborhood of the nodes, their diam = $O(\epsilon)$



Residual with support in $O(\epsilon)$ neighborhood of the nodes.

3. For every node O_i we get the Stokes eq. in dilated variables $\xi = x/\epsilon$



in an unbounded domain
independent of ϵ

4. These problems are
related by a problem
for the pressure on the
graph.

5. Equations on the graph.

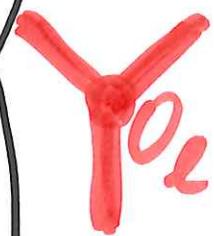
5.1 Steady state.

Find $p \in H^1(\Omega)$

$$-\frac{\partial^2}{\partial x_1^{(e)} \partial x_1^{(e)}} p(x_1^{(e)}) = 0, \text{ (linearity)}$$

 $x_1^{(e)} \in (0, l_e) \setminus e$

$$-\sum_{e: 0_l \in e} \alpha^{(e)} \frac{\partial p}{\partial x_1^{(e)}}(0) = 0, \text{ (flux conservation)}$$



$\forall 0_l, l=1, \dots, N_1$

$$\underline{\alpha}^{(e)} - \alpha^{(e)} \frac{\partial p}{\partial x_1^{(e)}}(0) = \sum_{i=N_1+1, \dots, N} \vec{g}_i \cdot \vec{n} ds$$

$$\underline{\alpha}^{(e)} = \int_{\tilde{\sigma}^{(e)}} U_p dx - \text{coef. prop.}$$

in the law flux-pressure slope.

For higher order approximations:
given pressure jumps in the nodes.

5.2. Time dependent equation on the graph

Operator $L^{(e)}$ relating the pressure drop and the flux in an infinite tube



Given $S \in L^2_{\text{per}}(0, T)$ find

$$v \in L^2_{\text{per}}(0, T; H_0^1(\sigma^{(e)})),$$

$$\frac{\partial v}{\partial \tau} \in L^2_{\text{per}}(0, T; L^2(\sigma^{(e)})) \quad \text{s.t.}$$

$$\begin{cases} \frac{\partial v}{\partial \tau}(y^{(e)}, \tau) - v \Delta' V(y^{(e)}, \tau) = S(\tau), \\ y^{(e)} \in \sigma^{(e)}, \tau \in \mathbb{R}, \end{cases} \quad (2)$$

$$V(y^{(e)}, \tau)|_{\partial \sigma^{(e)}} = 0$$

denote

$$L^{(e)} S(\tau) = \int_{\sigma^{(e)}} V(y^{(e)}, \tau) dy^{(e)}.$$

$$L^{(e)} : L^2_{\text{per}}(0, T) \longleftrightarrow H^1_{\text{per}}(0, T)$$

Problem on the graph

Given $\psi_\ell \in H^1_{\text{per}}(0, T)$, $\ell = 1, \dots, N$,

$F^{(e_i)} \in H^1_{\text{per}}(0, T; L^2(\mathcal{B}))$, $i = 1, \dots, M$

find $p \in L^2_{\text{per}}(0, T; H^1(\mathcal{B}))$ s.t.

$$\left. \begin{aligned} & \frac{\partial}{\partial x_1^{(e)}} \left(L^{(e)} \frac{\partial p}{\partial x_1^{(e)}}(x_1^{(e)}, \tau) \right) = F^{(e_i)}(x_1^{(e)}, \tau), \\ & x_1^{(e)} \in (0, |e|), \\ & \forall e = e_j, j = 1, \dots, M \\ & - \sum_{e: 0_e \in e} \left(L^{(e)} \frac{\partial p}{\partial x_1^{(e)}} \right)(0, \tau) = \psi_\ell(\tau), \\ & \ell = 1, \dots, N_1, \end{aligned} \right\} \quad (3)$$

$$- L^{(e)} \frac{\partial p}{\partial x_1^{(e)}}(0, \tau) = \psi_\ell(\tau), \quad \ell = N_1 + 1, \dots, N$$

$$\boxed{\sum_{i=1}^M \int_0^{l(e_i)} F^{(e_i)} dx_1^{(e_i)} + \sum_{\ell=1}^N \psi_\ell(\tau) = 0} \quad (iv)$$

compatibility cond.

Th.1. Let $H \in H^1_{per}(0, T)$.

$\exists!$ (V, S) satisfying (2) and

$$\int_{\sigma^{(e)}} V(y^{(e)}, \tau) dy^{(e)} = H(\tau)$$

$$\|V\|_{L^2(0, T; H_0^1(\sigma^{(e)}))} + \left\| \frac{\partial V}{\partial \tau} \right\|_{L^2(0, T; L^2(\sigma^{(e)}))}$$

$$+ \|S\|_{L^2(0, T)} \leq C \|H\|_{H^1(0, T)}$$

and $\exists C_e > 0$:

$$\forall Q \in L^2_{per}(0, T)$$

$$C_e^{-1} \|Q\|_{L^2(0, T)} \leq \|L^{(e)} Q\|_{H^1(0, T)} \leq$$

$$\leq C_e \|Q\|_{L^2(0, T)} \quad (\text{Galdi})$$

Th 2. Let $\langle \cdot \rangle = \frac{1}{T} \int_0^T \cdot d\tau$.

$$\langle L^e S \rangle = \mathcal{R}^e \langle S \rangle,$$

$$\alpha^{(e)} = \int_{\sigma^{(e)}} V(y') dy' > 0, V:$$

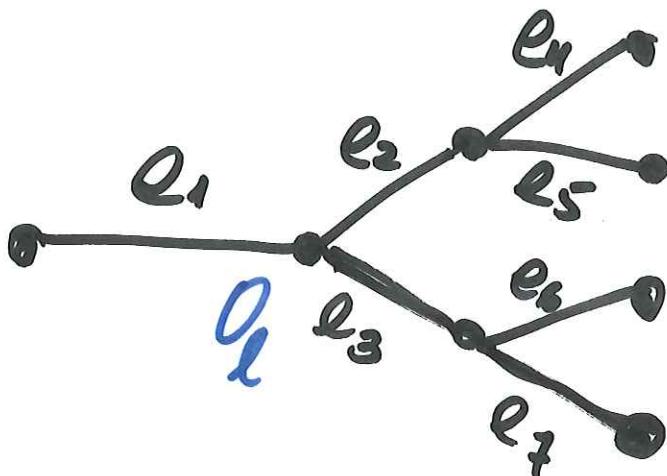
$$-y \Delta V = 1, y' \in \sigma^{(e)}, V|_{\partial \sigma^{(e)}} = 0.$$

Th. 3. Pb. (3) admits a unique
(up to an additive function of ε
from $L^2_{\text{per}}(0, T)$) solution p .

Proof. Lax-Milgram argument.

Conclusion : new Reynolds' type
equation on the graph,
non-local in time.

6. "Randomization"



p linear

$$-\sum_{e: \emptyset_i \in e} x^{(e)} \frac{\partial P}{\partial x_i^{(e)}}(0) = 0, \quad \forall i=1, \dots, N$$

$$P|_{e_i}^{(0)} = P|_{e_j}^{(0)} + \Delta P_{ij}, \quad i \neq j, \quad \emptyset_i \in e_i, e_j$$

$$-\frac{\partial P}{\partial x_i^{(e)}}(0) = \int_{\gamma_e^i} \vec{g}_i \cdot \vec{n} ds$$

ΔP_{ij} random variables (indep.)

Eg. for expectation; deviations?

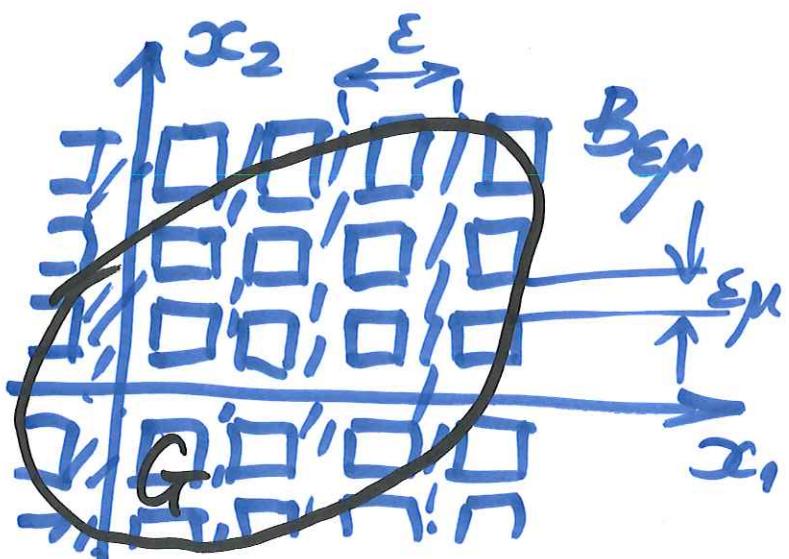
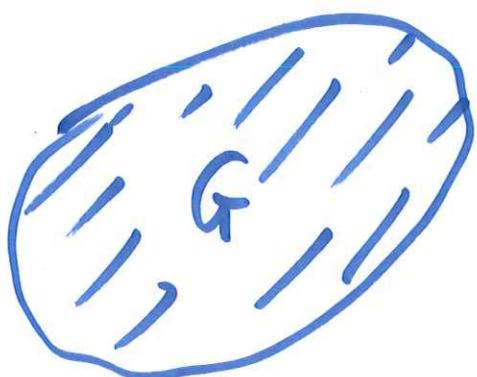
- Random sections $\sigma^{(i)}$
- Fractal extension :
distribution of the fluxes
in the edges of rank N
 $N \rightarrow +\infty$.
- Example of a random
thin structure :
Homogenization of a
random lattice structure

G.P. Comp. Math. Math. Phys. (USSR)
1983, 23, N5.

$G \subset \mathbb{R}^2$ bounded domain, $\forall G \in \mathcal{C}^2$

$$B_{\varepsilon\mu} = \bigcup_{k=-\infty}^{+\infty} \left\{ (x_1, x_2) \in \mathbb{R}^2 \mid |x_2 - k\varepsilon| < \frac{\varepsilon\mu}{2} \right\}$$

$$\bigcup \left\{ (x_1, x_2) \in \mathbb{R}^2 \mid |x_1 - k\varepsilon| < \frac{\varepsilon\mu}{2} \right\}$$



$$B_k^j = \left\{ (x_1, x_2) \in \mathbb{R}^2 \mid |x_{3-j} - k\varepsilon| < \frac{\varepsilon\mu}{2} \right\}, j=1,2,$$

\downarrow
 $A_{B_k^j}$ random valued constant 2×2 -matrix
 indep. in aggregate

$$P\{A_{B_k^j} = A^{(s)}\} = p_s, \quad s=1, \dots, 2$$

$$\sum_{s=1}^2 p_s = 1, \quad A^{(s)} = A^{(s)T} > 0, \quad A^{(s)} = (a_{ij}^{(s)})_{1 \leq i,j \leq 2}$$

$$\Pi = \bigcup_{k,j=-\infty}^{\infty} (B_k^1 \cap B_j^2) : A^{(0)} = A^{(0)T} > 0.$$

$$\left\{ \begin{array}{l} -\operatorname{div}(A \nabla u_{\varepsilon\mu}) = f(x), \quad x \in B_{\varepsilon\mu} \cap G, \\ A \nabla u_{\varepsilon\mu} \cdot \vec{n} = 0, \quad x \in \partial B_{\varepsilon\mu} \cap G, \\ u_{\varepsilon\mu} = 0, \quad x \in \partial G \cap \bar{B}_{\varepsilon\mu}, \end{array} \right.$$

where $f \in C^1(\bar{G})$,

$$A = \begin{cases} A_{B_k^j} & \text{in } B_k^j \setminus \Pi, \\ A^{(0)} & \text{in } \Pi. \end{cases}$$

Denote $\bar{A} = (\bar{a}_{ij})_{1 \leq i, j \leq 2}$,

$$\bar{a}_{11} = \frac{\sum_{s=1}^2 p_s}{2} \frac{a_{11}^{(s)} - a_{12}^{(s)} (a_{22}^{(s)})^{-1} a_{21}^{(s)}}{2}, \quad \bar{a}_{12} = 0,$$

$$\bar{a}_{22} = \frac{\sum_{s=1}^2 p_s}{2} \frac{a_{22}^{(s)} - a_{21}^{(s)} (a_{11}^{(s)})^{-1} a_{12}^{(s)}}{2}, \quad \bar{a}_{21} = 0,$$

$$u_0 : \begin{cases} -\operatorname{div}(\bar{A} \nabla u_0) = f(x), \quad x \in G, \\ u_0 |_{\partial G} = 0 \end{cases}$$

Th. $\forall \delta \in (0, 1)$

$$P \left\{ \frac{\|u_{\varepsilon\mu} - u_0\|_{L^2(B_{\varepsilon\mu} \cap G)}}{\sqrt{\operatorname{mes}(B_{\varepsilon\mu} \cap G)}} > (\sqrt{\varepsilon} + \sqrt{\mu})^{1-\delta} \right\} \leq C \varepsilon^\delta.$$