

Quantitative particle approximation of nonlinear Fokker-Planck equations with singular kernels

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joint work with C. Olivera (Campinas) and A. Richard (CentraleSupélec) to appear in Ann Sc Norm Super Pisa'21+

Workshop METASTABILITY, MEAN-FIELD PARTICLE SYSTEMS
AND NON LINEAR PROCESSES

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Overview

Introduction

Rate of convergence to the PDE

The nonlinear process and propagation of chaos

Recent progress and perspectives

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Rate of convergence to the PDE

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Recent progress and perspectives

Nonlinear Fokker-Planck equation with singular interaction

- ▶ We study on $[0, T) \times \mathbb{R}^d$

$$\begin{cases} \partial_t u(t, x) = \Delta u(t, x) - \nabla \cdot (u(t, x) K *_x u(t, x)), \\ u(0, x) = u_0(x), \end{cases} \quad (\text{NLFP})$$

with K **locally integrable** and **singular** at 0.

- ▶ Our main interest : stochastic particle approximation of (NLFP).
- ▶ Why?
 - ▶ Macroscopic to microscopic description (and back!);
 - ▶ Numerical schemes...

Classical approach: mean-field interactions

- ▶ (NLFP) is seen as the FP equation for the **non-linear process**

$$\begin{cases} dX_t = K * u_t(X_t)dt + \sqrt{2}dW_t, \\ \mathcal{L}(X_t) = u_t. \end{cases} \quad (\text{NLSDE})$$

- ▶ Particle system in **mean-field interaction** :

$$dX_t^{i,N} = \frac{1}{N} \sum_{j=1}^N K(X_t^{i,N} - X_t^{j,N}) + \sqrt{2}dW_t^{i,N} \quad (\text{PS})$$

and its empirical measure $\mu_t^N = \frac{1}{N} \sum_{i=1}^N \delta_{X_t^{i,N}}$.

- ▶ **Propagation of chaos** : convergence in law of μ_t^N in $\mathcal{P}(\mathcal{C})$ towards the law of X .
- ▶ Issues: **K singular** \rightarrow wellposedness of (PS), (NLSDE) and the propagation of chaos?
- ▶ Probabilistic approach to **singular** interactions:
 - ▶ Boltzmann, Burgers, Navier-Stokes, Keller-Segel equations, ...
 - ▶ studied by Bossy, Calderoni, Cattiaux, Fournier, Hauray, Jabin, Jourdain, Méléard, Mischler, Osada, Pulvirenti, Talay, ...

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Another viewpoint: moderate interaction

Motivated by singular attractive kernels for which either:

- ▶ existence of (PS) is unknown ;
- ▶ or existence ok, but convergence unknown.

Moderately interacting particles :

$$dX_t^{i,N} = F\left(V^N * (K * \mu_t^N)(X_t^{i,N})\right)dt + \sqrt{2}dW_t^{i,N},$$

where:

- ▶ $V^N(x) = N^{d\alpha}V(N^\alpha x)$, $\alpha > 0$ and V a regular density;
- ▶ F a smooth cut-off.
- ▶ Some references : [Oelschläger'85], [Méléard & Roelly'87], [Jourdain & Méléard'98]

→ A **semigroup approach** was recently developed by Flandoli, Olivera and their collaborators to get uniform (non-quantitative) convergence of $V^N * \mu^N$ towards a mild solution to:

FKPP, 2d Navier-Stokes equations, PDE-ODE system related to aggregation phenomena.

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Our main objectives

What assumptions on K and what suitable functional framework for (NLFP) so the following holds?

- ▶ Convergence of $\{\mu_t^N = \frac{1}{N} \sum_{i=1}^N \delta_{X_t^{i,N}}, t \in [0, T]\}$ to the solution (NLFP) when $N \rightarrow \infty$:
 - ▶ which range of α ?
 - ▶ what is the rate of convergence ?
- ▶ Well-posedness of (NLSDE).
- ▶ Propagation of chaos towards (NLSDE) (without the cut-off and the mollifier)

Which kernels?

A typical example in dimension $d \geq 2$ is the family of **Riesz kernels**:

$$K_s(x) = \pm \nabla \mathcal{V}_s(x)$$

where

$$\mathcal{V}_s(x) := \begin{cases} |x|^{-s} & \text{if } s \in (0, d-1) \\ -\log|x| & \text{if } s = 0 \end{cases}, \quad x \in \mathbb{R}^d.$$

Examples:

- ▶ Coulomb interactions: K_s , with $s = d - 2$ ($d \geq 3$);
- ▶ 2d Navier-Stokes equation (vorticity): $K(x) = \frac{x^\perp}{|x|^2}$;
- ▶ Parabolic-elliptic Keller-Segel model: $K(x) = -\chi \frac{x}{|x|^d}$ (*attractive...*);
- ▶ Some attractive-repulsive kernels.

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Assumptions on K and α

(H_K) :

- (i) $K \in L^p(\mathcal{B}_1)$, for some $p \in [1, +\infty)$;
- (ii) $K \in L^q(\mathcal{B}_1^c)$, for some $q \in [1, +\infty)$;
- (iii) There exists $r \geq \max(p', q')$ and $\zeta \in (0, 1]$ such that:

$$\mathcal{N}_\zeta(K * f) \lesssim \|f\|_{L^1 \cap L^r(\mathbb{R}^d)}, \quad \forall f \in L^1 \cap L^r(\mathbb{R}^d).$$

(\mathcal{N}_ζ the Hölder seminorm of parameter ζ .)

(H_α) : α and r satisfy

$$0 < \alpha < \frac{1}{d + 2d(\frac{1}{2} - \frac{1}{r}) \vee 0}.$$

► Notice here that for $f \in L^1 \cap L^r(\mathbb{R}^d)$ one has

$$\|K * f\|_{L^\infty(\mathbb{R}^d)} \leq C_{K,d} \|f\|_{L^1 \cap L^r(\mathbb{R}^d)}.$$

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Convergence of the mollified empirical measure

Assume (always today) that $u_0 \geq 0$ and $\int_{\mathbb{R}^d} u_0 = 1$.

Proposition

Let K satisfying (H_K) (i)-(ii), $u_0 \in L^1 \cap L^r(\mathbb{R}^d)$ with $r \geq \max(p', q')$.
There exists $T > 0$ and a unique $u : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}$ s.t.

$$u \in \mathcal{C}([0, T]; L^1 \cap L^r(\mathbb{R}^d))$$

and

$$u_t = e^{t\Delta} u_0 - \int_0^t \nabla \cdot (e^{(t-s)\Delta} (u_s K * u_s)) ds, \quad 0 \leq t \leq T.$$

Denote by T_{max} the maximal existence time.

Convergence of the mollified empirical measure

Theorem ([O.-R.-T. Ann Sc Norm Super Pisa'21+])

Let $T < T_{max}$ and assume (H_K) and (H_α) . Under suitable conditions on the initial conditions, we have for $\{u_t^N = V^N * \mu_t^N, t \in [0, T]\}_{N \in \mathbb{N}}$:
 $\forall \varepsilon > 0$ and $\forall m \geq 1$,

$$\left\| \|u^N - u\|_{T, L^1 \cap L^r(\mathbb{R}^d)} \right\|_{L^m(\Omega)} \lesssim \left\| \|u_0^N - u_0\|_{L^1 \cap L^r(\mathbb{R}^d)} \right\|_{L^m(\Omega)} + N^{-\varrho + \varepsilon},$$

where

$$\varrho = \min \left(\alpha \zeta, \frac{1}{2} \left(1 - \alpha \left(d + d \left(1 - \frac{2}{r} \right) \vee 0 \right) \right) \right).$$

Convergence of the mollified empirical measure

Without the cutoff F in the drift of the particles, we get:

Corollary

For any $\varepsilon \in (0, \varrho)$, any $\eta > 0$ and any $m \geq 1$,

$$\mathbb{P} \left(\|u_t^N - u_t\|_{T, L^1 \cap L^r(\mathbb{R}^d)} \geq \eta \right) \lesssim \frac{1}{\eta^m} \left(\left\| \|u_0^N - u_0\|_{L^1 \cap L^r(\mathbb{R}^d)} \right\|_{L^m(\Omega)} + N^{-\varrho + \varepsilon} \right)^m.$$

Some consequences and remarks

- ▶ Same rate for the **genuine empirical measure** of (PS)

$$\left\| \sup_{t \in [0, T]} \|\mu_t^N - u_t\|_0 \right\|_{L^m(\Omega)} \leq C N^{-\varrho + \varepsilon},$$

where $\|\cdot\|_0$ denotes the Kantorovich-Rubinstein metric

- ▶ The rate in the previous results holds almost surely.
- ▶ Cannot expect here a $N^{-\frac{1}{2}}$ rate of convergence because of the short range interactions : “best possible” $N^{-\alpha}$.

- ▶ **Coulomb-type kernels** (like Biot-Savart in $d = 2$, Riesz with $s = d - 2$),
 - ▶ convergence happens for $\alpha < \frac{1}{2(d-1)}$ ($d = 2 \Rightarrow \alpha = (\frac{1}{2})^-$);
 - ▶ best possible rate is $\varrho = \left(\frac{1}{2(d+1)}\right)^-$ (obtained for $\alpha = \left(\frac{1}{2(d+1)}\right)^+$, $r = +\infty$, $\zeta = 1$).
- ▶ **Keller-Segel parabolic elliptic** ($d = 2$: global solution $\chi < 8\pi$, blow up in finite time otherwise).
 - ▶ we get the above rate for any value of χ ;
 - ▶ the result holds even if the PDE explodes in finite time ($\chi > 8\pi$).
- ▶ **The Riesz kernels with $s > d - 2$** do not satisfy (H_K) (iii). However, by imposing more regularity on the initial conditions and smaller values of α , we get a rate of convergence for singular Riesz kernels with $s \in (d - 2, d - 1)$.

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Sketch of proof

- ▶ Derive the SPDE satisfied by the mollified empirical measure u^N in its mild form

$$u_t^N(x) = e^{t\Delta} u_0^N(x) - \int_0^t \nabla \cdot e^{(t-s)\Delta} \langle \mu_s^N, V^N(x - \cdot) F(K * u_s^N(\cdot)) \rangle ds \\ - \frac{1}{N} \sum_{i=1}^N \int_0^t e^{(t-s)\Delta} \nabla V^N(x - X_s^{i,N}) \cdot dW_s^i, \quad x \in \mathbb{R}^d$$

applying Itô's formula to $G_{t,V^N}(s, x - \cdot) := e^{(t-s)\Delta} V^N(x - \cdot)$ on each particle between 0 and t for x and t fixed. (sum up, $\frac{1}{N}$, rearrange..)

- ▶ For $q \geq 1$ establish that

$$\sup_{N \in \mathbb{N}^*} \mathbb{E} \left[\sup_{t \in [0, T]} \|u_t^N\|_{L^q(\mathbb{R}^d)}^q \right] < \infty.$$

(using the above mild form and combining Gronwall lemma and "martingale" estimate)

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Decompose $\|u_t^N - u_t\|_{L^1 \cap L^r}$ into several terms (u is the unique sln of the mild equation with cut-off):

$$u_t^N(x) - u_t(x) = e^{t\Delta}(u_0^N - u_0)(x) + E_t(x) - M_t^N(x) + \int_0^t \nabla \cdot e^{(t-s)\Delta} \left(u_s F(K * u_s) - u_s^N F(K * u_s^N) \right) (x) ds,$$

where we have set

$$E_t(x) := \int_0^t \nabla \cdot e^{(t-s)\Delta} \langle \mu_s^N, V^N(x - \cdot) \left(F(K * u_s^N(x)) - F(K * u_s^N(\cdot)) \right) \rangle ds,$$

$$M_t^N(x) := \frac{1}{N} \sum_{i=1}^N \int_0^t e^{(t-s)\Delta} \nabla V^N(x - X_s^{i,N}) \cdot dW_s^i. \quad (1)$$

Pivoting in the last term and using that

$$\|K * f\|_{L^\infty(\mathbb{R}^d)} \leq C_{K,d} \|f\|_{L^1 \cap L^r(\mathbb{R}^d)}$$

$$\begin{aligned} \|u_t^N - u_t\|_{L^1 \cap L^r(\mathbb{R}^d)} &\leq \|u_0^N - u_0\|_{L^1 \cap L^r(\mathbb{R}^d)} + C \int_0^t \frac{1}{\sqrt{t-s}} \|u_s^N - u_s\|_{L^1 \cap L^r(\mathbb{R}^d)} ds \\ &\quad + \|E_t\|_{L^1 \cap L^r(\mathbb{R}^d)} + \|M_t^N\|_{L^1 \cap L^r(\mathbb{R}^d)}. \end{aligned}$$

Using the Grönwall lemma for convolution integrals, we obtain

$$\begin{aligned} \|u^N - u\|_{t, L^1 \cap L^r(\mathbb{R}^d)} &\leq C \left(\|u_0^N - u_0\|_{L^1 \cap L^r(\mathbb{R}^d)} \right. \\ &\quad \left. + \|E\|_{t, L^1 \cap L^r(\mathbb{R}^d)} + \|M^N\|_{t, L^1 \cap L^r(\mathbb{R}^d)} \right). \end{aligned} \tag{2}$$

It remains to control the moments of $\|E\|_{t, L^q(\mathbb{R}^d)}$ and $\|M\|_{t, L^q(\mathbb{R}^d)}$ for $q \in \{1, r\}$. This is where the two expressions in the definition of the speed ρ appear.

Observe that (positivity of F , heat kernel)

$$\|E_t\|_{L^q(\mathbb{R}^d)} \leq C \int_0^t \frac{1}{(t-s)^{\frac{1}{2}}} \left(\int_{\mathbb{R}^d} \langle \mu_s^N, V^N(x-\cdot) \right. \\ \left. \left| F(K * u_s^N(\cdot)) - F(K * u_s^N(x)) \right|^q dx \right)^{\frac{1}{q}} ds.$$

F Lipschitz and the ζ -Hölder continuity of $K * u^N$ give

$$\|E_t\|_{L^q(\mathbb{R}^d)} \leq C \int_0^t \frac{\|u_s^N\|_{L^1 \cap L^r(\mathbb{R}^d)}}{(t-s)^{\frac{1}{2}}} \left(\int_{\mathbb{R}^d} \langle \mu_s^N, V^N(x-\cdot) |\cdot-x|^\zeta \rangle^q dx \right)^{\frac{1}{q}} ds.$$

Since V is compactly supported (wlog, assume $\text{supp}(V) \subset B_1$), we have that $V^N(x-y) |y-x|^\zeta \leq N^{-\alpha\zeta} V^N(x-y)$. Thus

$$\|E_t\|_{L^q(\mathbb{R}^d)} \leq \frac{C}{N^{\alpha\zeta}} \int_0^t \frac{1}{(t-s)^{\frac{1}{2}}} \|u_s^N\|_{L^1 \cap L^r(\mathbb{R}^d)} \|u_s^N\|_{L^q(\mathbb{R}^d)} ds \\ \leq \frac{C}{N^{\alpha\zeta}} \int_0^t \frac{1}{(t-s)^{\frac{1}{2}}} \|u_s^N\|_{L^1 \cap L^r(\mathbb{R}^d)}^2 ds,$$

Apply Hölder's inequality with $p = \frac{3}{2}$ to obtain

$$\|E_t\|_{L^q(\mathbb{R}^d)} \leq \frac{C}{N^{\alpha\zeta}} \left(\int_0^t (t-s)^{-\frac{3}{4}} ds \right)^{\frac{2}{3}} \left(\int_0^t \|u_s^N\|_{L^1 \cap L^r(\mathbb{R}^d)}^6 ds \right)^{\frac{1}{3}}.$$

Finally, we have from Jensen's inequality that for $m \geq 3$,

$$\left\| \|E\|_{t, L^q(\mathbb{R}^d)} \right\|_{L^m(\Omega)} \leq \frac{C}{N^{\alpha\zeta}} \left(\int_0^t \mathbb{E} \left[\|u_s^N\|_{L^1 \cap L^r(\mathbb{R}^d)}^{2m} \right] ds \right)^{\frac{1}{m}}$$

and the bound uniform bounds on u^N permits to conclude that

$$\left\| \|E\|_{t, L^q(\mathbb{R}^d)} \right\|_{L^m(\Omega)} \leq \frac{C}{N^{\alpha\zeta}}. \quad (3)$$

This inequality immediately extends to $1 \leq m < 3$.

Sketch of proof

Main issue: control the moments of

$$\sup_{t \leq T} \left\| \frac{1}{N} \sum_{i=1}^N \int_0^t e^{(t-s)\Delta} \nabla V^N(X_s^i - \cdot) dW_s^i \right\|_{L^1 \cap L^r(\mathbb{R}^d)}.$$

- ▶ Not a martingale, fix the time in the heat operator and it becomes one;
- ▶ to control the $L^1 \cap L^r(\mathbb{R}^d)$ norm, use stochastic integration techniques in infinite-dimensional spaces [van Neerven et al.'07];
- ▶ use Garsia-Rodemich-Rumsey's lemma to put the sup inside.

Note that this is where the limitation on $\alpha(H_\alpha)$ arises.

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Definition

Assume (H_K) (i)-(ii), $u_0 \in L^1 \cap L^r(\mathbb{R}^d)$ with $r \geq \max(p', q')$ and consider the canonical space $\mathcal{C}([0, T]; \mathbb{R}^d)$ equipped with its canonical filtration. We say that $\mathbb{Q} \in \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))$ solves the **nonlinear martingale problem** associated to (NLSDE) if:

- (i) $\mathbb{Q}_0 = u_0$;
- (ii) For any $t \in (0, T]$, \mathbb{Q}_t has density q_t and $q \in \mathcal{C}([0, T]; L^1 \cap L^r(\mathbb{R}^d))$;
- (iii) For any $f \in \mathcal{C}_c^2(\mathbb{R}^d)$, the process $(M_t)_{t \in [0, T]}$ defined as

$$M_t := f(w_t) - f(w_0) - \int_0^t \left[\Delta f(w_s) + \nabla f(w_s) \cdot (K * q_s(w_s)) \right] ds$$

is a \mathbb{Q} -martingale, where $(w_t)_{t \in [0, T]}$ denotes the canonical process.

Proposition

Let $T < T_{max}$. Then the martingale problem related to (NLSDE) is well-posed.

⇒ the McKean-Vlasov equation (NLSDE) admits a unique weak solution \tilde{X} . Combined with the convergence theorem, it comes:

$$\left\| \max_{i \in \{1, \dots, N\}} \sup_{t \in [0, T]} |X_t^{i, N} - \tilde{X}_t^i| \right\|_{L^m(\Omega)} \leq C N^{-\varrho + \varepsilon}, \quad \forall N \in \mathbb{N}^*.$$

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Theorem

Same hypotheses as in previous Theorem + Assume that $\{X_0^i, i \in \mathbb{N}\}$ are identically distributed and that $\langle u_0^N, \varphi \rangle \xrightarrow{\mathbb{P}} \langle u_0, \varphi \rangle, \forall \varphi \in \mathcal{C}_b(\mathbb{R}^d)$.
Then

$$\mu^N \xrightarrow{(d)} \mathbb{Q},$$

where \mathbb{Q} is the law of the solution of (NLSDE).

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Sketch of proof

Usual strategy:

- (1) Consider the nonlinear MP with cutoff;
- (2) Prove the tightness of the family $\Pi^N := \mathcal{L}(\mu^N)$ in the space $\mathcal{P}(\mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d)))$;
- (3) Prove that any limit point Π^∞ of Π^N is $\delta_{\mathbb{Q}}$.
- (4) Lift the cutoff.

(3) is the most technical part:

- ▶ work in fact on

$$\mathcal{H} := \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d)) \times \mathcal{C}([0, T]; L^1 \cap L^r(\mathbb{R}^d))$$

with $\tilde{\Pi}^N = \mathcal{L}(\mu^N, u^N)$ (as in [Méléard & Roelly'87]).

- ▶ introduce a quadratic functional Γ on \mathcal{H} , which depends on the form of the martingale problem.
- ▶ use the convergence of $\tilde{\Pi}^N$ to $\tilde{\Pi}^\infty$ to prove that $\Gamma = 0$ $\tilde{\Pi}^\infty$ -a.e. This is where μ^N and the particle system appear.
- ▶ deduce that the first coordinate of $\tilde{\Pi}^\infty$ solves the nonlinear MP (with cutoff).

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Overview

Introduction

Rate of convergence to the PDE

The nonlinear process and propagation of chaos

Recent progress and perspectives

Recent progress

- ▶ In a recent work [Guo & Luo '21] extend our method to particles with common noise:

$$dX_t^{i,N} = V^\varepsilon * (K * \mu_t^N)(X_t^{i,N})dt + \sqrt{2\nu} \sum_k \sigma_k^N(X_t^{i,N})dW_t^k,$$

and quantify its convergence (in a two step procedure) to

$$\partial_t u(t, x) = \nu \Delta u(t, x) - \nabla \cdot (u(t, x) K *_x u(t, x))$$

for a class of kernels such as repulsive Riesz kernels for $s \in [0, d-2]$.

- ▶ Rate of convergence for the IPS without cutoff, working on the torus:

$$\limsup_{N \rightarrow +\infty} N^{\varrho-\varepsilon} \sup_{t \in [0, T]} \|u_t^N - u_t\|_{L^p(\mathbb{T}^d)} \leq X \text{ a.s.}$$

- ▶ Extension to Burgers.

Some next steps

- ▶ Numerical applications : use our result to quantify the convergence of a scheme coming from the moderately interacting particles.
- ▶ Treat non-Markovian particle systems : e.g. the parabolic-parabolic Keller-Segel model.
- ▶ Improve the constraint on α by changing the functional space.

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